# Molecular Force Fields of s-trans-1,3-Butadiene and the Second Stable Conformer

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Normal-coordinate analyses of s-trans-1,3-butadiene and its various deuterated and  $^{13}C$ -substituted analogs were performed using the group-coordinate force field, and a set of force constants was determined. The infrared spectra of the second stable conformer of 1,3-butadiene- $d_0$ , -1,1,4,4- $d_4$ , and - $d_6$  were analyzed on the basis of normal-coordinate calculations of the planar s-cis form and a nonplanar gauche form. Such an analysis suggests that the second stable conformer takes a gauche form rather than the s-cis. To increase the number of reliable data to be used in this study, the infrared and Raman spectra of the  $d_0$ , 1,1,4,4- $d_4$ , and  $d_6$  species were observed in low-temperature Ar matrices. The infrared spectra of the second stable conformer were also observed using a combination of the high-temperature nozzle and low-temperature matrix techniques.

1,3-Butadiene has been a target molecule of numerous structural, spectroscopic, and theoretical studies. Part of these studies has treated the structure of the most stable s-trans form of this molecule1-5) and the assignments of the vibrational spectra arising from this form. 6-15) On the other hand, various efforts have been made toward the elucidation of the internalrotation potential and the second stable conformer around the central C-C bond. 16-23) This subject has been studied also by various methods of quantumchemical calculations. 24-33) Interest in this direction of study has been renewed by the recent observation of the infrared bands associated with the second stable conformer trapped in low-temperature Ar matrices. 34-36) The in-plane vibrational frequencies as well as the inplane force constants of the s-trans and s-cis forms were calculated by an ab initio method37) and the normalcoordinate analyses based on the observed data for deuterated and <sup>13</sup>C-substituted species were also performed.36)

The purpose of the present study is twofold. (1) One is to obtain reliable force constants which can be conveniently used (at least as initial values) in the normal-coordinate analyses of longer molecules with conjugated double bonds, for example, isomers of 1,3,5hexatriene, poly(acetylene), various isomers of retinal Although many authors have and  $\beta$ -carotene, etc. already reported various sets of force constants of s-transand s-cis-1,3-butadiene, 13-15,36,37) those force constants do not seem satisfactory for such a purpose. (2) The other is to examine if it is possible to obtain information on the internal-rotation angle of the second stable conformer through the normal-coordinate analyses. Whether the second stable conformer is in the planar s-cis form or in the nonplanar (gauche) form is a longstanding problem which still remains to be solved.

Reliable experimental data are absolutely necessary for this kind of study. Therefore, we have observed the infrared and Raman spectra of 1,3-butadiene and its  $1,1,4,4-d_4$  and  $d_6$  analogs in low-temperature Ar matrices. Also, we have observed the infrared spectra of these compounds using a combination of the high-temperature nozzle and the low-temperature Ar

matrix techniques, in order to obtain the bands due to the second stable conformer. In this way we have confirmed the results of other authors for the  $d_0$  and  $d_6$ -samples<sup>34–36)</sup> and have obtained the data for the 1,1,4,4- $d_4$  analog for the first time.

### Method

Infrared and Raman Measurements. Infrared and Raman spectra were observed for 1,3-butadiene- $d_0$  and two deuterated analogs (1,1,4,4- $d_4$  and  $d_6$ ) in low-temperature Ar matrices. The  $d_0$  sample was obtained from Takachiho Kagaku Kogyo Co., Ltd. (purity 99%) and the  $d_4$  and  $d_6$  samples from Merck Sharp and Dohme Canada Ltd. (98 atom% D). No further purification was undertaken.

Infrared spectra were recorded on a Hitachi 260-50 infrared spectrophotometer. Raman spectra were observed on a Raman spectrometer consisting of a Spex 1401 double monochromator, an NEC GLG 3300 Arion laser, an HTV 649 photomultiplier and a home-made photon counter.

The low-temperature apparatus for the low-temperature matrix infrared measurements consisted of a CTI Model 21 Cryocooler and a sample chamber equipped with a CsI plate, KBr windows, and a fine Pyrex nozzle through which the Ar/sample mixture was passed. The temperature of the nozzle could be raised to about 770 K with a surrounding tungsten heater. The temperature of the Ar/sample mixture was monitored with an alumel-chromel thermocouple placed inside the nozzle. For low-temperature Raman measurements an Osaka Oxygen Co. Cryo Mini D equipped with a cold head of silver-coated copper and a sample chamber with quartz windows was used.

Table 1. Internal coordinates

	TABLE 1. INTERNAL COORDINATES
Symbol	Description
$\Delta r_{ij}$	Stretching of the bond linking the atoms
	i and j
$\Delta oldsymbol{\phi_{ijk}}$	Bending of the angle between the two bonds
	$r_{ij}$ and $r_{ik}$
$\Delta  heta_{ ext{i-jkl}}$	Out-of-plane bending of the plane formed
	by the atoms i(apex), j, k, and l.
$\Delta t_{ij}$	Torsion around the bond $r_{ij}$

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(b) 
$$H^{1}$$
  $H^{4}$   $H^{8}$   $H^{9}$   $H^{6}$   $H^{1}$   $H^{4}$   $H^{6}$   $H^{1}$   $H^{2}$   $H^{2}$   $H^{3}$   $H^{4}$   $H^{6}$   $H^{8}$   $H^{8}$ 

Fig. 1. Numbering of the atoms of (a) s-trans- and (b) s-cis-1,3-butadiene.

Normal Coordinate Analysis. The normal coordinate analyses of 1,3-butadiene- $d_0$  and its deuterated and <sup>13</sup>C-substituted analogs were performed using the program NCTB38) on a FACOM M190 computer system at the Laboratory of International Collaboration on Elementary Particle Physics, Faculty of Science, the University of Tokyo.

The structure and numbering of atoms of s-trans- and s-cis-1,3-butadiene are shown in Fig. 1. The notations for the internal coordinates are given in Table 1. The out-of-plane bending coordinate  $(\Delta \theta_{i-jkl})$  and the torsional coordinate  $(\Delta t_{ij})$  are defined according to the IUPAC recommendations.<sup>39)</sup> Therefore,  $\Delta \theta_{i-ikl}$ is defined as

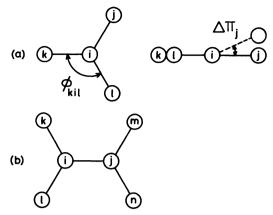


Fig. 2. Numbering of the atoms and the angles relating to the definitions of (a) out-of-plane bending and (b) torsion.

$$\Delta heta_{i-jkl} = \Delta \pi_{j} \sin \phi_{kil}$$

where  $\Delta \pi_i$  and  $\phi_{kil}$  are given in Fig. 2(a). The sign of  $\Delta\theta_{i-jkl}$  is defined in such a way that  $\Delta\pi_i$  is plus when the atom j moves above the plane defined by the atoms k, i, and l. The torsional coordinate  $\Delta t_{ij}$  is defined as

$$\Delta t_{ij} = (\Delta \tau_{kijm} + \Delta \tau_{kijn} + \Delta \tau_{lijm} + \Delta \tau_{lijn})/4,$$

where  $\Delta \tau_{xijy}$  denotes the change in the individual torsional angle between the bonds xi and yj [see Fig. 2(b)].

Group coordinates were defined on the basis of the internal coordinates as given in Table 2. (Previously the group coordinates were called the local-symmetry coordinates.40)

In the first step of the normal coordinate analysis the

Table 2.	. Group coordinates <sup>a)</sup>
Group coordinate	Mode [Abbreviation]
$S_{\mathbf{a}}^{1} = (\Delta r_{12} + \Delta r_{23}) / \sqrt{2}$	$CH_{\frac{1}{2}}$ symmetric stretching $[\nu_{s}(CH_{\frac{1}{2}})]$
$S_{a}^{2} = (\Delta r_{67} + \Delta r_{78})/\sqrt{2}$	$CH_2^2$ symmetric stretching $[\nu_s(CH_2^2)]$
$S_{b}^{1} = (\Delta r_{12} - \Delta r_{23}) / \sqrt{2}$	$CH_2^1$ antisymmetric stretching $[\nu_a(CH_2^1)]$
$S_b^2 = (\Delta r_{67} - \Delta r_{78}) / \sqrt{2}$	$\mathrm{CH_2^2}$ antisymmetric stretching $[ u_{\mathrm{a}}(\mathrm{CH_2^2})]$
$S_{ m c}^{ m 1} = \Delta r_{ m 45}$	$\mathrm{CH^1}$ stretching $[\nu(\mathrm{CH^1})]$
$S_{\mathrm{c}}^{2} = \Delta r_{\mathrm{90}}$	$ m CH^2$ stretching [ $ u(CH^2)$ ]
$S_{ extbf{d}}^{ extbf{1}} = \Delta r_{25}$	$C=C^1$ stretching $[\nu(C-C^1)]$
$S_{\mathbf{d}}^{2} = \Delta r_{70}$	$ ext{C=C}^2$ stretching $\left[ u( ext{C=C}^2) ight]$
$S_{\rm e}^1 = (2\Delta\phi_{123} - \Delta\phi_{325} - \Delta\phi_{521})/\sqrt{6}$	$ ext{CH}_2^1  ext{ scissors } [\delta( ext{CH}_2^1)]$
$S_{\rm e}^2 = (2\Delta\phi_{\rm 678} - \Delta\phi_{\rm 870} - \Delta\phi_{\rm 076})/\sqrt{6}$	$ ext{CH}_2^2  ext{ scissors } [\delta( ext{CH}_2^2)]$
$S_{ m f}^1 = (\Delta \phi_{325} - \Delta \phi_{521})/\sqrt{2}$	$CH_2^1$ rocking $[\gamma_r(CH_2^1)]$
$S_{\rm f}^2 = (\Delta \phi_{870} - \Delta \phi_{076})/\sqrt{2}$	$CH_2^2$ rocking $[\gamma_r(CH_2^2)]$
$S_{\mathbf{g}}^{1} = (2\Delta\phi_{250} - \Delta\phi_{054} - \Delta\phi_{452})/\sqrt{6}$	$C=C-C^1$ deformation $[\delta(CCC^1)]$
$S_{\rm g}^2 = (2\Delta\phi_{705} - \Delta\phi_{509} - \Delta\phi_{907})/\sqrt{6}$	C=C-C <sup>2</sup> deformation $[\delta(CCC^2)]$
$S_{\rm h}^1 = (\Delta \phi_{452} - \Delta \phi_{054}) / \sqrt{2}$	$CH^1$ in-plane bending $[\delta(CH^1)]$
$S_{\rm h}^2 = (\Delta \phi_{907} - \Delta \phi_{509})/\sqrt{2}$	$CH^2$ in-plane bending $[\delta(CH^2)]$
$S_{\mathbf{i}} = \Delta r_{50}$	$C-C$ stretching $[\nu(C-C)]$
$S_1^1 = \Delta \theta_{2-135}$	$CH_2^1$ wagging $[\gamma_w(CH_2^1)]$
$S_{\mathbf{j}}^{\mathbf{z}} = \Delta \theta_{\mathbf{7-680}}$	$\text{CH}_2^2 \text{ wagging } [\gamma_{\mathbf{w}}(\text{CH}_2^2)]$
$S_{\mathbf{k}}^{\mathbf{i}} = \Delta \theta_{\mathbf{5-420}}$	$CH^1$ out-of-plane bending $[\gamma_w(CH^1)]$
$S_{\mathbf{k}}^{2} = \Delta \theta_{0-975}$	$CH^2$ out-of-plane bending $[\gamma_w(CH^2)]$
$S_1^1 = \Delta t_{25}$	$CH_2^1$ twisting $[\gamma_t(CH_2^1)]$
$S_1^2 = \Delta t_{70}$	$CH_2^2$ twisting $[\gamma_t(CH_2^2)]$
$S_{\mathbf{m}} = \Delta t_{50}$	C–C torsion $[\tau(C-C)]$

a) The index zero in the subscripts corresponds to the atom 10.

Tai	BLE 3. DEFINITION OF GROUP-COORDINATE					
	FORCE CONSTANTS					
Symbol	Description					
(1) In-pla	ne force constants					
(i) Diag	gonal force constants					
$F_\mathtt{a}$	CH <sub>2</sub> symmetric stretching					
$F_{\mathbf{a}}'$	CD <sub>2</sub> symmetric stretching					
$F_{ m a}^{\prime\prime}$	CHD symmetric stretching					
$F_{\mathtt{b}}$	CH <sub>2</sub> antisymmetric stretching					
$F_{\mathbf{b}}'$	CD <sub>2</sub> antisymmetric stretching					
$F_{ m b}^{\prime\prime}$	CHD antisymmetric stretching					
$F_{\mathbf{c}}$	CH stretching					
$F_{c}^{\prime}$	CD stretching					
$F_{\mathbf{d}}$	C=C stretching					
$F_{ m e}$	CH <sub>2</sub> scissors					
$F_{\mathbf{f}}$	CH <sub>2</sub> rocking					
$\dot{F_{\mathbf{g}}}$	C=C-C deformation					
$F_{ m h}$	CH in-plane bending					
$F_{i}^{-}$	C-C stretching					
	diagonal force constants					
$F_{\mathtt{dd}}$	C=C stretching/C=C stretching					
$F_{ m de}$	C=C stretching/CH <sub>2</sub> scissors					
$F_{ t dg}$	C=C stretching/C=C-C deformation					
$F_{ m dh}$	C=C stretching/CH in-plane bending					
$F_{\mathtt{di}}$	C=C stretching/C-C stretching					
$F_{ m eh}$	CH <sub>2</sub> scissors/CH in-plane bending					
$F_{ m fh}$	CH <sub>2</sub> rocking/CH in-plane bending					
$F_{xi}$	C=C-C deformation/C-C stretching					
$F_{ m hi}$	CH in-plane bending/C-C stretching					
$F_{gg}$	C=C-C deformation/C=C-C deformation					
$F_{ m hh}$	CH in-plane bending/CH in-plane bending					
(2) Out-of	-plane force constants					
(i) Diago	onal force constants					
$F_{\mathbf{j}}$	CH <sub>2</sub> wagging					
$F_{\mathbf{k}}$	CH out-of-plane bending					
$F_{\scriptscriptstyle  m I}$	CH <sub>2</sub> twisting					
$F_{ m m}$	C–C torsion					
(ii) Off-diagonal force constants						
$F_{\mathtt{jk}}$	CH <sub>2</sub> wagging/CH out-of-plane bending					
$F_{\mathbf{k}\mathbf{l}}$	CH out-of-plane bending/CH <sub>2</sub> twisting					
	$(\underline{CH}_2 = \underline{CH} -)$					
$F_{\mathtt{k}\mathtt{k}}$	CH out-of-plane bending/CH out-of-plane					
	bending					
$F_{\mathrm{kl}}'$	CH out-of-plane bending/CH <sub>2</sub> twisting					
_	$(\underline{CH}_2 = \underline{CH} - \underline{CH} = \underline{CH}_2)$					

modified Urey-Bradley-Shimanouchi force field40) based on the work of Abe13) was assumed. However, it became clear that some force constants were strongly correlated and accordingly had large uncertainties. Then, we employed the group-coordinate force field (GCFF),41) which was previously called the localsymmetry force field (LSFF).40) The definition of the GC force constants used are given in Table 3.

CH, twisting/CH, twisting

 $F_{11}$ 

s-trans-1,3-Butadiene: The planar C<sub>2h</sub> structure<sup>4)</sup> with  $r_{\text{C=C}} = 1.342 \text{ Å}, r_{\text{C-C}} = 1.463 \text{ Å}, r_{\text{C-H}} = 1.093 \text{ Å}, \phi_{\text{C=C-C}} = 1.093 \text{ Å}$ 123.6°, and  $\phi_{C=C-H}$  (three kinds)=120.9° was used in the present calculation. The least-squares adjustments of the in-plane GC force constants were made against 96 observed frequencies in total (11 for  $d_0$ , 11 for 1,1,4,4-

 $d_4$ , 11 for  $d_6$ , 9 for 1,1,2- $d_3$ , 9 for trans, trans-1,4- $d_2$ , 8 for  $cis, cis-1, 4-d_2$ , 7 for 2,3-d<sub>2</sub>, 10 for trans-1-d<sub>1</sub>, 4 for  $cis-1-d_1$ , 7 for  $1,4^{-13}C_2$ , and 9 for  $1^{-13}C_1$ ). The observed frequencies of the  $d_0$ , 1,1,4,4- $d_4$ , and  $d_6$  samples were obtained from the present low-temperature matrix infrared and Raman measurements. The observed frequencies of the other deuterated and 13C-substituted analogs were selected from Refs. 13-15 and 36. The following assumptions were made to obtain a good fit between the calculated and observed frequencies using a relatively small number of force constants (9 diagonal and 11 off-diagonal). (a) The CH2 and CH stretching frequencies (CD2, CD, and CHD stretching as well) were not included in the least-squares adjustments because of their uncertainties due to possible Fermi resonances. The diagonal force constants ( $F_a$ and  $F_b$  for  $CH_2$  and  $F_c$  for CH) were fixed to the values given in Table 4 and no off-diagonal constants were taken into account. Slightly different values were used for the diagonal force constants of  $CD_2$  ( $F'_a$  and  $F'_b$ ), CD ( $F'_c$ ), and CHD ( $F''_a$  and  $F''_b$ ). (b) The off-diagonal force constants  $F_{\rm di}({\rm C=C/C-C})~{\rm and}~F_{\rm dd}({\rm C=C/C=C})~{\rm had}$ large uncertainties when they were included in the leastsquares adjustments. Therefore, these force constants were fixed to the values which were obtained from an ab initio calculation.<sup>37)</sup>

For the out-of-plane vibrations nine force constants listed in Tables 3 and 4 were determined by the leastsquares method against 74 observed frequencies in total (7 for  $d_0$ , 7 for 1,1,4,4- $d_4$ , 7 for  $d_6$ , 6 for 1,1,2- $d_3$ , 5 for trans, trans-1,4- $d_2$ , 5 for cis, cis-1,4- $d_2$ , 5 for cis, trans-1,4 $d_2$ , 4 for 2,3- $d_2$ , 6 for trans-1- $d_1$ , 5 for cis-1- $d_1$ , 6 for 2- $d_1$ , 6 for  $1,4^{-13}C_2$ , and 5 for  $1^{-13}C_1$ ). Here again, the observed data of the  $d_0$ , 1,1,4,4- $d_4$ , and  $d_6$  samples were from the present measurements, while those of the other deuterated and 13C-substituted analogs were taken from Refs. 13-15 and 36. Choice of the kinds of force constants used was somewhat arbitrary, but the agreement between the calculated and observed frequencies was quite satisfactory.

Second Stable Conformer of 1,3-Butadiene: Since the structure of the second stable conformer of 1,3-butadiene has not been experimentally determined, two model structures, s-cis and gauche, were chosen for normal coordinate analysis. For the s-cis form all the bond lengths and bond angles were set equal to those of the s-trans form and the internal-rotation angle around the central  $\mathrm{C_{5}\text{--}C_{10}}$  bond was rotated by  $180^{\circ}$  from s-trans as shown in Fig. 1(b). The same kinds of force constants as used for s-trans were taken into account. Normalcoordinate calculations were carried out in two steps. At first the force constants of s-trans were transferred to s-cis and the observed infrared bands of the latter were assigned with reference to the results of this calculation. The force constants of s-cis were then adjusted in order to obtain the best fit between the calculated and observed frequencies. In this procedure the force constants associated with the in-plane vibrations of the terminal CH<sub>2</sub> group were fixed to the values of s-trans. The offdiagonal force constants,  $F_{dd}$  and  $F_{di}$ , were again fixed to the values obtained from an ab initio calculation. 37) The other six force constants,  $F_d$ ,  $F_g$ ,  $F_h$ ,  $F_i$ ,  $F_{gg}$ , and

Table 4. Group-coordinate force constants of *s-trans*, *s-cis*, and *gauche* 1,3-butadiene and their uncertainties (in parentheses). Units are mdyn Å $^{-1}$  (=10 $^{2}$  N m $^{-1}$ ) for stretching and stretch-stretch interaction force constants, mdyn (=10 $^{-8}$  N) for stretch-bend interaction force constants, and mdyn Å (=10 $^{-18}$  N m) for bending, bend-bend interaction, and torsional force constants.

Force constant	s-trans	s-cis	gauche
In-plane			
$F_{\mathbf{a}}$	5.068	5.068	5.068
$F_{\mathbf{a}}'$	5.214	5.214	5.214
$F_{ m a}^{\prime\prime}$	5.570		•
$F_{\mathtt{b}}$	5.057	5.057	5.057
$F_{ m b}'$	5.123	5.123	5.123
$F_{\mathrm{b}}^{\prime\prime}$	4.724		<del>_</del>
$^{\cdot}F_{\mathbf{c}}$	4.995	4.995	4.995
$F_{\mathbf{c}}'$	4.981	4.981	4.981
$F_{\sf d}$	8.886 (0.105)	8.976(0.122)	9.003 (0.075)
$ar{F_{e}}$	0.455 (0.003)	0.455	0.455
$ec{F_{f}}$	0.577 (0.007)	0.577	0.577
$F_{\mathbf{g}}$	0.691 (0.011)	0.699 (0.193)	0.546 (0.070)
$oldsymbol{F_{\mathtt{h}}}$	0.519(0.006)	0.532(0.011)	0.531(0.007)
$F_{\mathbf{i}}^{"}$	5.428 (0.164)	4.964 (0.286)	4.914(0.155)
$F_{ m dd}$	-0.1	-0.1	-0.1
$F_{ t de}$	-0.197(0.012)	-0.197	-0.197
$F_{ t dg}$	0.257 (0.095)	0.257	0.257
$F_{ m dh}$	0.232(0.030)	0.232	0.232
$F_{ m di}$	0.4	0.4	0.4
$F_{ m eh}$	-0.013(0.004)	-0.013	-0.013
$F_{ m fh}$	0.056 (0.005)	0.056	0.056
$F_{\mathbf{gi}}$	0.082 (0.049)	0.082	0.082
$F_{ m hi}$	-0.225(0.023)	-0.225	-0.225
$F_{gg}^{}$	0.073 (0.010)	0.008 (0.193)	-0.026(0.071)
$F_{ m hh}^{ m ss}$	0.037 (0.004)	0.024(0.011)	0.021(0.007)
Out-of-plane	` ,	,	, ,
$F_{f j}$	0.297 (0.001)	0.297	0.297
$oldsymbol{f_k}$	0.409 (0.005)	0.352 (0.006)	0.387 (0.015)
$F_1^{"}$	0.511(0.005)	0.533 (0.009)	0.550(0 018)
$\dot{F_{ m m}}$	0.109 (0.005)	0.046(0.003)	0.055 (0.009)
$F_{\mathtt{jk}}^{ ext{ in}}$	0.055 (0.002)	0.045 (0.003)	0.038 (0.008)
$F_{\mathbf{k}\mathbf{l}}$	-0.027(0.002)	-0.010(0.003)	-0.028 (0.009)
$F_{\mathbf{k}\mathbf{k}}$	-0.005(0.004)	0.024 (0.005)	0.013 (0.015)
$F_{\mathbf{k}\mathbf{l}}^{'}$	0.023 (0 002)	-0 026(0 003)	-0.014(0.009)
$F_{11}^{"}$	-0.017(0.005)	0.011 (0.008)	-0.003(0.017)

 $F_{\rm hh}$ , were adjusted freely by the least-squares method using 30 observed frequencies in total (6 for  $d_0$ , 7 for  $1,1,4,4-d_4$ , 7 for  $d_6$ , 5 for  $1,4^{-13}C_2$ , 5 for  $1^{-13}C_1$ ). The observed data for the  $1,4^{-13}C_2$  and  $1^{-13}C_1$  analogs were taken from Ref. 36. As for the out-of-plane vibrations, the eight force constants,  $F_{\rm k}$ ,  $F_{\rm l}$ ,  $F_{\rm m}$ ,  $F_{\rm jk}$ ,  $F_{\rm kl}$ ,  $F_{\rm kk}$ ,  $F_{\rm kl}$ , and  $F_{\rm ll}$  were adjusted freely by the least-squares method using 28 observed frequencies (7 for  $d_0$ , 6 for  $1,1,4,4-d_4$ , 4 for  $d_6$ , 5 for  $1,4^{-13}C_2$ , 6 for  $1^{-13}C_1$ ). Only the CH<sub>2</sub> wagging force constant,  $F_{\rm j}$ , was fixed to the value of *s-trans*.

As the structural parameters of the gauche form the values calculated by Skaarup et al.<sup>30)</sup> were slightly modified with reference to the deviations between the ab initio and experimental results for the s-trans form. The parameters used were as follows:  $r_{\text{C=C}}=1.341$  Å,  $r_{\text{C-C}}=1.474$  Å,  $r_{\text{C-H}}=1.092$  Å,  $\phi_{\text{C=C-C}}=119.3^{\circ}$ ,  $\phi_{\text{C=C-H}}(5-2-1)=121.5^{\circ}$ ,  $\phi_{\text{C=C-H}}(5-2-1)=121.8^{\circ}$ , and  $t(\text{internal-rotation angle around } C_5-C_{10})=121.8^{\circ}$ 

40° from s-cis. In the normal-coordinate analysis of the gauche form no interaction force constants between the "in-plane" and "out-of-plane" coordinates were taken into account. Under the same assumptions as made for the s-cis case the 14 force constants were determined from 58 observed frequencies by the least-squares method.

## Results

The infrared and Raman spectra of 1,3-butadiene and its  $1,1,4,4-d_4$  and  $d_6$  analogs observed in low-temperature Ar matrices are shown in Figs. 3 and 4. Since the Ar/sample ratio was not high, the sample molecules might not be completely isolated in the Ar matrix. However, the sharpness of individual bands inherent to the low-temperature matrix measurements made it possible to detect many weak bands. The Raman spectra of these compounds in low-temperature

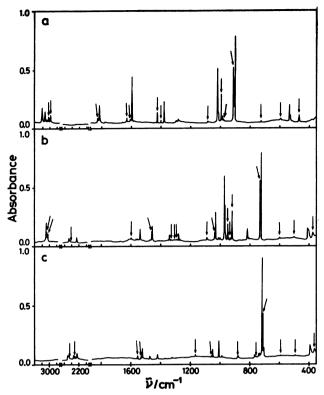


Fig. 3. Infrared spectra of (a) 1,3-butadiene-d<sub>0</sub>, (b) -1,1,4,4-d<sub>4</sub>, and (c) -d<sub>6</sub> observed in Ar matrices at 20 K. Mixtures of Ar and the samples (Ar/sample ratio 250) were heated to about 770 K in the nozzle and then deposited onto the cold CsI plate. The arrows indicate the bands due to the second stable conformer.

Ar matrices were observed in this study for the first time. The infrared spectra in Fig. 3 were observed from the low-temperature Ar matrices obtained by depositing a mixture of Ar and the sample which had been heated to about 770 K. Therefore, the intensities of the bands due to the second stable conformer are enhanced in Fig. 3 and they are indicated with arrows. The bands without arrows are associated with the most stable s-trans form. The bands with arrows were found to decrease in intensity on irradiation of light from a Xe-lamp, indicating that these bands were due to the second stable conformer.35) Our present results for the  $d_0$  and  $d_6$  samples (Figs. 3a and 3c) agreed with the findings of previous authors, 34-36) except for some Huber-Wälchli and Günthard<sup>36)</sup> minor features. observed more bands than we did; for example, we could not observe the medium-intensity band at 1740 cm<sup>-1</sup> of the  $d_0$  species which they assigned to a combination tone of the second stable conformer. The infrared spectrum of the second stable conformer of the 1,1,4,4d<sub>4</sub> species (Fig. 3b) was obtained for the first time.

The results of normal-coordinate analyses will be described below.

#### **Discussion**

Band Assignments and Force Constants of the s-trans Form. As described earlier, we started the normal-coordinate analyses of 1,3-butadiene and its deuterated and <sup>13</sup>C-

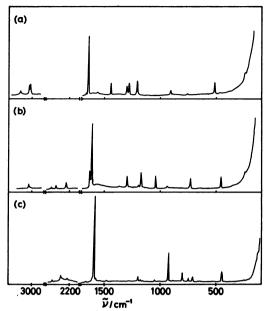


Fig. 4. Raman spectra of (a) 1,3-butadiene- $d_0$ , (b) -1,1,4,4- $d_4$ , and (c) - $d_6$  observed in Ar matrices (Ar/sample ratio 50) at 15 K. Excitation wavelength, 514.5 nm.

Table 5. Observed and calculated vibrational frequencies of s-trans-1,3-butadiene- $d_0$  and mode assignments

Species	No.	Obsd v/cm <sup>-1</sup>		Mode and PED <sup>a</sup> )
a	ν,	3105	3102	ν <sub>•</sub> (CH <sub>2</sub> *) 49
g	$\nu_{\rm g}$	3025		$\nu$ (CH*) 47
	$\nu_3$	3014		$\nu_{8}(CH_{2}^{*})$ 46
	$\nu_{4}$	1644		$\nu(C=C^*)$ 36, $\nu(C-C)$ 21
	$\nu_5$	1441	1446	
	$\nu_{\rm s}$	1279	1296	
	$\nu_{7}$	1206	1203	, , , , ,
	•			$\gamma_r(CH_2^*)$ 11
	$\nu_{8}$	887	901	$\nu(C-C)$ 37, $\gamma_r(CH_2^*)$ 29
	ν	513	510	
$\mathbf{a_u}$	$\nu_{10}$	1022	1019	$\gamma_{\rm t}({\rm CH_2^*})$ 31, $\gamma_{\rm w}({\rm CH^*})$ 18
_	$\nu_{11}$	905	908	$\gamma_{\mathbf{w}}(\mathbf{CH_{2}^{*}})$ 51
	$\nu_{12}$	535	527	$\gamma_{\rm w}({\rm CH^*})$ 28, $\gamma_{\rm t}({\rm CH_2^*})$ 21
	$\nu_{13}$	163 <sup>b)</sup>	162	τ(C–C) 89
$\mathbf{b_g}$	$\nu_{14}$	974	969	$\gamma_{\rm w}({\rm CH^*}) 30, \gamma_{\rm t}({\rm CH_2^*}) 25$
	$\nu_{15}$	908	910	$\gamma_{\mathbf{w}}(\mathbf{CH_2^*})$ 50
	$\nu_{16}$	754	751	$\gamma_{\rm t}({\rm CH_2^*})$ 24, $\gamma_{\rm w}({\rm CH^*})$ 18
$\mathbf{b_u}$	$\nu_{17}$	3103	3101	$\nu_a(CH_2^*)$ 50
	$\nu_{18}$	3062	3029	ν(CH*) 46
	$ u_{19}$	2986	2997	$\nu_{\rm s}({\rm C=}{\rm H_2^*})~48$
	$\nu_{20}$	1597	1601	$\nu(C=C^*)$ 37, $\delta(CH_2^*)$ 17
	$\nu_{21}$	1381	1387	$\delta(\text{CH}_2^*)$ 33, $\nu(\text{C=C*})$ 10
	$\nu_{22}$	1297	1293	$\delta(\text{CH*})$ 33
	$ u_{23}$	988	995	$\gamma_{\rm r}({\rm CH_2^*}) \ 40, \ \delta({\rm CH^*}) \ 16$
	$\nu_{24}$	301b)	300	$\delta(CCC^*)$ 55

a) The asterisk in the mode notation indicates that the two group coordinates of the same mode (having superscript 1 and 2) have equal PED's; for example,  $\nu_a(\text{CH}_2^*)$  49 represents  $\nu_a(\text{CH}_2^1)$  49 and  $\nu_a(\text{CH}_2^2)$  49. In this Table PED's are not normalized to a total of 100. b) Ref. 18.

Table 6. Observed and calculated vibrational frequencies of s-trans-1,3-butadiene-1,1,4,4-d<sub>4</sub>

And mode assignments

Species	No.	Obsd $\tilde{v}/\mathrm{cm}^{-1}$	Calcd $\tilde{\nu}/\text{cm}^{-1}$	Mode and PEDa)
a <sub>g</sub>	$\nu_1$	3026	3038	ν(CH*) 49
•	$\nu_2$	2316	2340	$\nu_{a}(CD_{2}^{*})$ 49
	$\nu_3^-$	2225	2233	$\nu_{s}(CD_{2}^{*})$ 46
	$\nu_4$	1612	1602	$\nu(C=C^*)$ 34, $\nu(C-C)$ 26
	$\nu_5$	1298	1296	$\delta(\text{CH*})$ 30
	$\nu_{6}$	1170	1183	ν(C–C) 46
	$\nu_7$	1040	1036	$\delta(\mathrm{CD_2^*})$ 40
	$\nu_8$	740	726	$\gamma_{\rm r}({\rm CD_2^*})$ 34
	$\nu_{9}$	454	462	$\delta(CCC^*)$ 31
$\mathbf{a}_{\mathbf{u}}$	$\nu_{10}$	968	957	$\gamma_{\rm w}({\rm CH^*})$ 25, $({\rm CD_2^*})$ 23
	$\nu_{11}$	722	724	$\gamma_{\mathbf{w}}(\mathbf{CD_2^*})$ 51
	$\nu_{12}$	404	399	$\gamma_{\rm t}({\rm CD_2^*})$ 29, $\gamma_{\rm w}({\rm CH^*})$ 2
	$\nu_{13}$	149 <sup>b)</sup>	148	τ(C–C) 89
$\mathbf{b_g}$	$\nu_{14}$	940	934	$\gamma_{\rm w}({\rm CH^*})$ 43
	$\nu_{15}$	728		$\gamma_{\mathbf{w}}(\mathrm{CD_2^*})$ 44
	$\nu_{16}$	610°)	602	$\gamma_{\mathbf{t}}(\mathbf{CD_{2}^{*}})$ 31
$\mathbf{b_u}$	$\nu_{17}$	3037		ν(CH*) 50
	$\nu_{18}$	2337	2339	$\nu_a(\mathrm{CD_2^*})$ 49
	$\nu_{19}$	2228		$\nu_s({ m CD_2^*})$ 46
	$\nu_{20}$	1533		ν(C=C*) 44
	$\nu_{21}$	1274	1263	$\delta(\text{CH*})$ 45
	$\nu_{22}$	1028	1022	· • ·
	$\nu_{23}$	811	811	$\gamma_{\mathbf{r}}(\mathrm{CD_2^*})$ 44
	$\nu_{24}$	258°)	262	$\delta(CCC^*)$ 52

a) See the footnote a) of Table 5. b) Ref. 18. c) Ref. 13.

Table 7. Observed and calculated vibrational frequencies of s-trans-1,3-butadiene- $d_6$  and mode assignments

	AND MODE ASSIGNMENTS					
Species	No.	Obsd v/cm <sup>-1</sup>	Calcd $\tilde{\nu}/\mathrm{cm}^{-1}$	Mode and PED <sup>a</sup> )		
ag	$\nu_1$	2349	2342	$\nu_{a}(\mathrm{CD_{2}^{*}})$ 47		
	$\nu_2$	2275	2275	$\nu({\rm CD*}) 34,  \nu_{\rm s}({\rm CD_2^*})  10$		
	$\nu_3$	2218	2220			
	$\nu_4$	1592	1586	ν(C=C*) 36, ν(C-C) 29		
	$\nu_5$	1199	1190	$\nu(\text{CC})$ 37, $\delta(\text{CD}_2^*)$ 11		
	$\nu_6$	1050	1030			
	$\nu_7$	925	918	$\delta(CD^*)$ 35		
	$\nu_8$	743	723	$\gamma_{\rm r}({\rm CD_2^*})$ 32, $\nu({\rm C-C})$ 24		
	$\nu_9$	445	451	$\delta(CCC^*)$ 31		
$\mathbf{a}_{\mathrm{u}}$	$\nu_{10}$	766	759	$\gamma_{\rm w}({\rm CD}^*)$ 23, $\gamma_{\rm t}({\rm CD}_2^*)$ 22		
	$\nu_{11}$	715	717	$\gamma_{\mathbf{w}}(\mathbf{CD_2^*})$ 41		
	$\nu_{12}$	391	389	$\gamma_{\mathbf{w}}(\mathbf{CD^*})$ 25, $\gamma_{\mathbf{t}}(\mathbf{CD_2^*})$ 24		
	$\nu_{13}$	142 <sup>b)</sup>	- 139	τ(C-C) 92		
b <sub>e</sub>	$\nu_{14}$	800	796	$\gamma_{w}(CD^{*})$ 43, $\gamma_{w}(CD_{2}^{*})$ 14		
•	$\nu_{15}$	705	702			
	$\nu_{16}$	603°)	592	$\gamma_{\mathbf{t}}(\mathbf{CD_2^*})$ 39		
$\mathbf{b_u}$	$\nu_{17}$	2350	2339	$\nu_{a}(CD_{2}^{*})$ 49		
_	$\nu_{18}$	2269	2263	$\nu(CD) 23, \nu_s(CD_2^*) 22$		
	ν <sub>19</sub>	2222	2209	$\nu(CD) 25, \nu_{s}(CD_{2}^{*}) 25$		
	$\nu_{20}$	1521	1511	ν(C=C*) 43		
	$\nu_{21}$	1046	1043	$\delta(\mathrm{CD_2^*})$ 30, $\delta(\mathrm{CD^*})$ 14		
	$\nu_{22}$	1006	989	$\delta(\mathrm{CD_2^*})$ 16, $\delta(\mathrm{CD^*})$ 15,		
				$\gamma_r(CD_2^*)$ 10		
	$\nu_{23}$	738	738			
	$\nu_{24}$	250	255	$\delta(CCC^*)$ 53		

a) See the footnote a) of Table 5. b) Ref. 18. c) Ref. 13.

Table 8. Observed and calculated vibrational frequencies of s-cis-1,3-butadiene- $d_0$  and mode assignments

Species	No.	Obsd $\tilde{\nu}/\mathrm{cm}^{-1}$	$_{\tilde{\nu}/\mathrm{cm}^{-1}}^{\mathrm{Calcd}}$	Mode and PEDa)
a <sub>1</sub>	$\nu_1$	3103	3101	ν <sub>a</sub> (CH <sub>2</sub> *) 49
	$\nu_2$	3014	3033	ν(CH*) 47
	$\nu_3$	2986	2997	$\nu_{\rm s}({ m CH_2^*}) \ 47$
	$\nu_4$	1633	1627	$\nu(C=C^*)$ 38, $\nu(C-C)$ 17
	$\nu_5$	1425	1432	$δ(CH_2^*)$ 34, $ν(C-C)$ 15
	$\nu_{6}$		1338	$\delta(CH^*)$ 27
	$\nu_7$	1087	1075	$\gamma_{\rm r}({\rm CH_2^*})$ 22, $\nu({\rm C-C})$ 16,
				$\delta(\text{CH*})$ 13
	$\nu_8$		878	$\nu(C-C)$ 54, $\gamma_r(CH_2^*)$ 20
	$\nu_9$		256	δ(CCC*) 47
$\mathbf{a_2}$	$\nu_{10}$	9835)	980	$\gamma_{\rm t}({\rm CH_2^*})$ 29, $\gamma_{\rm w}({\rm CH^*})$ 22
_	$\nu_{11}$	915°)	914	$ν_{\mathbf{w}}(\mathbf{CH_2^*})$ 50
	$ u_{12}$	727 <sup>b)</sup>	729	$\gamma_{\rm w}({\rm CH}^*)$ 23, $\gamma_{\rm t}({\rm CH}_2^*)$ 19
	$\nu_{13}$	136 <sup>d</sup> )	135	τ(C–C) 97
$b_1$	$\nu_{14}$	3103	3101	$\nu_{\rm a}({ m CH_2^*}) \ 49$
	$\nu_{15}$	3014	3034	ν(CH*) 47
	$\nu_{16}$	2986	2998	$\nu_{\rm s}({ m CH_2^*}) \ 47$
	$\nu_{17}$	1612	1623	$\nu(C=C^*)$ 37, $\delta(CH_2^*)$ 12
	$\nu_{18}$	1403	1415	$\delta(\text{CH}_2^*)$ 35, $\delta(\text{CH}^*)$ 11
	$\nu_{19}$		1278	$\delta(CH^*)$ 23
	$\nu_{20}$		1043	$\gamma_{\rm r}({\rm CH_2^*})$ 35, $\delta({\rm CH^*})$ 13
	$\nu_{21}$	596	591	$\delta(CCC^*)$ 42
$\mathbf{b_2}$	$\nu_{22}$	996	997	$\gamma_{\rm t}({\rm CH_2^*})$ 36, $\gamma_{\rm w}({\rm CH^*})$ 14
	$\nu_{23}$	914	909	$\gamma_{\mathbf{w}}(\mathbf{CH_{2}^{*}})$ 51
	$\nu_{24}$	470	470	$\gamma_{\rm w}({\rm CH^*}) 40, \gamma_{\rm t}({\rm CH_2^*}) 15$

a) See the footnote a) of Table 5. b) See the text as for the assignment of these bands. c) Estimated from  $\nu_{11}+\nu_{23}=1829~{\rm cm}^{-1}$ . d) Ref. 21.

substituted analogs using the modified Urey-Bradley-Shimanouchi force constants obtained by Abe. 13) Since the band assignments by Abe and other authors<sup>36)</sup> mostly seemed reasonable, we refined the GC force constants on the basis of their assignments. By fixing the force constants relating to the carbon-hydrogen stretching  $(F_a, F_b, \text{ and } F_c)$  and the interaction force constants between the skeletal stretchings  $(F_{dd} \text{ and } F_{di})$ to reasonable values (see the section of Method), a convergent set of force constants was obtained with satisfactorily small uncertainties as given in Table 4. The results of calculation for the  $d_0$ , 1,1,4,4- $d_4$ , and  $d_6$ species in the s-trans form are shown in Tables 5-7, together with the observed frequencies. The results of calculation for the other deuterated and <sup>13</sup>C-substituted analogs are available on request.

Internal-rotation Angle of the Second Stable Conformer. The normal frequencies of the planar s-cis and nonplanar gauche forms were calculated in the same manner as is described in the section on Method. The results of calculations for the  $d_0$  species are given in Tables 8 and 9. Similar tables for the  $1,1,4,4-d_4$  and  $d_6$  species are available from the authors on request. The force constants obtained for the s-cis and gauche forms are also listed in Table 4. Although most of the observed frequencies can be interpreted on the basis of either the s-cis or gauche form, the presence of a very weak band at

Table 9. Observed and calculated vibrational frequencies of gauche-1,3-butadiene- $d_0$  and mode assignments

Species	No.	Obsd ṽ/cm <sup>-1</sup>	Calcd v/cm <sup>-1</sup>	Mode and PED <sup>a</sup> )		
a	$\nu_1$	3103	3098	$\nu_{\rm a}({ m CH_2^*})$ 50		
	$\boldsymbol{\nu_2}$	3014	3032	ν(CH*) 46		
	$\nu_3$	2986	3000	ν <sub>s</sub> (CH <sub>2</sub> *) 47		
	$\nu_4$	1633	1629	ν(C=C*) 38, (C–C) 18		
	$ u_5$	1425	1432	δ(CH <sub>2</sub> *) 36, $ν$ (C–C) 13		
	$\nu_6$		1328	δ(CH*) 29		
	$\nu_7$	1087	1081	$\gamma_{\rm r}({\rm CH_2^*})$ 24, $\nu({\rm C-C})$ 16,		
				$\delta(\text{CH*})$ 12		
	$\nu_8$	983	982	$\gamma_{\rm t}({\rm CH_2^*})$ 29, $\gamma_{\rm w}({\rm CH^*})$ 23		
	$\nu_9$	915 <sup>b)</sup>	915	$\gamma_{\mathbf{w}}(\mathbf{CH_2^*})$ 47		
	$\nu_{10}$		872	$\nu(\text{CC})$ 55, $\gamma_{r}(\text{CH}_{2}^{*})$ 18		
	$\nu_{11}$	727	731	$\gamma_{\rm w}({\rm CH^*}) \ 21, \ \gamma_{\rm t}({\rm CH_2^*}) \ 18$		
	$\nu_{12}$		256	$\delta(CCC^*)$ 46		
	$\nu_{13}$	136°)	137	τ(C–C) 89		
b	$\nu_{14}$	3103	3098	$\nu_{\rm a}({ m CH_2^*})$ 50		
	$\nu_{15}$	3014	3032	ν(CH*) 46		
	$\nu_{16}$	2986	3001	$\nu_{\rm s}({ m CH_2^*}) \ 47$		
	$\nu_{17}$	1612	1619	$\nu(\text{C=C*}) \ 37, \ \delta(\text{CH}_2^*) \ 13$		
	$\nu_{18}$	1403	1413	$\delta(\mathrm{CH_2^*})$ 34, $\delta(\mathrm{CH^*})$ 11		
	$\nu_{19}$		1278	δ(CH*) 24		
	$\nu_{20}$		1040	$\gamma_{\rm r}({ m CH_2^*}) \ 36, \delta({ m CH^*}) \ 11$		
	$\nu_{21}$	996	996	$\gamma_{\rm t}({\rm CH_2^*})$ 35, $\gamma_{\rm w}({\rm CH^*})$ 15		
	$ u_{22}$	914	907	$\gamma_{\mathbf{w}}(\mathbf{CH_2^*})$ 50		
	$ u_{23}$	596	597	$\delta(\text{CCC*})$ 20, $\gamma_{\text{w}}(\text{CH*})$ 17		
	$\nu_{24}$	470	465	$\delta(\text{CCC*})$ 23, $\gamma_{\mathbf{w}}(\text{CH*})$ 19		

a) See the footnote a) of Table 5. b) Estimated from  $\nu_9 + \nu_{22} = 1829$  cm<sup>-1</sup>. c) Ref. 21.

983 cm<sup>-1</sup> and a weak band at 727 cm<sup>-1</sup> is better acounted for by assuming the gauche form. When the s-cis form is assumed, both of these bands must be assigned to the infrared-inactive  $a_2$  modes,  $\nu_{10}$  and  $\nu_{12}$ , in which the CH<sub>2</sub> twisting and CH out-of-plane bending are almost equally coupled. The corresponding bands are found at 927 cm<sup>-1</sup> ( $\nu_{10}$ ) and 594 cm<sup>-1</sup> ( $\nu_{12}$ ) for the 1,1,4,4-d<sub>4</sub> species and at 587 cm<sup>-1</sup> ( $\nu_{12}$ ) for the d<sub>6</sub> species. It is likely that the  $a_2$  modes of the s-cis form (C<sub>2v</sub>) are made infrared-active due to the lowering of molecular symmetry from C<sub>2v</sub> (s-cis) to C<sub>2</sub> (gauche).

As described above, the second stable conformer probably exists in the gauche form rather than in the s-cis form. However, it is interesting to note that most of the infrared bands due to the second stable conformer can be reasonably assigned on assuming the planar s-cis form. Therefore, as far as the vibrational assignment for a less stable conformer of a compound with conjugated double bonds is concerned, it seems justifiable to use the s-cis form as a practical model instead of paying too much attention to the problem of internal-rotation angle which is inherently difficult to determine.

Assignments of Infrared and Raman Bands. Most of the infrared and Raman bands shown in Figs. 3 and 4 could be assigned. The tables which list the observed infrared and Raman frequencies together with their assignments are available from the authors on request.

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